Direct numerical simulations of injection flows using cylindrical and spherical coordinates

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> Paolo Orlandi: *A vortical and turbulent life*, Roma, September 20, 2014.



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Conclusion

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3D numerical codes: coordinates system

I. Danaila's code (CYL) University of Rouen (cylindrical coordinates)



B. J. Boersma's code (SPH) TU Delft, The Netherlands (spherical coordinates)



Outline

Navier-Stokes numerical codes using cylindrical and spherical coordinates

Utility of the DNS of axisymmetric injection flows

- Investigation of the physics of vortex rings
- Improve models for the injection velocity profile
- Application in automotive industry
- New nice mathematical developments
- From fluids to superfluids

4 Conclusion



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3D numerical codes: equations

- Rai & Moin, JCP, 1991.
- Verzicco & Orlandi, JCP, 1996.
- Orlandi, Kluwer Academic Press, 1999.

Cylindrical coordinates: variables $(q_{\theta} = v_{\theta}, q_r = v_r \cdot r, q_z = v_z, p).$

Navier-Stokes equations

- incompressible $(\operatorname{div} \vec{v} = 0)$
- low Mach number approximation $(M \rightarrow 0)$



Discretization

Mesh : 3D, staggered

- θ and z directions: uniform grid
- *r* direction: stretched grid (cyl coordinates).



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Spatial discretization: second order finite differences. Time integration:

- (CYL) convective terms: 3-steps Runge-Kutta method,
- (CYL) diffusive terms: Crank-Nicolson method,
- (SPH) explicit Adams-Bashfort scheme.



Conclusion

Fractional time step (projection) method

for each step of Runge-Kutta:

momentum equations (ADI factorization)

$$\left(1 - \frac{\alpha_l}{2}\Delta t \mathcal{A}_c\right)\Delta \hat{q}_c^{\prime} = \left[\gamma_l \mathcal{H}_c^{\prime} + \rho_l \mathcal{H}_c^{\prime-1} - \alpha_l \mathcal{G}_c p^{\prime} + \alpha_l \mathcal{A}_c q_c^{\prime}\right]$$

• Poisson equation(FFT in θ + cyclic reduction)

$$\mathcal{L}\Phi^{l+1} = \frac{1}{\alpha_l \Delta t} \, \mathcal{D}\hat{\hat{q}}^{\,l}$$

corrected velocity field

$$q^{l+1} = -\alpha_l \Delta t \mathcal{G} \Phi + \hat{q}^l$$

scalar equation (TVD scheme)



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Navier-Stokes equations for low Mach flows

Idea: remove the pressure waves ($\epsilon = \gamma M^2$)

$$\frac{\partial \rho_0}{\partial t} + \nabla . (\rho_0 \vec{v}_0) = \mathbf{0}$$

$$\frac{\partial \rho_0 \vec{v}_0}{\partial t} + \nabla . \left(\rho_0 \ \vec{v}_0 \otimes \vec{v}_0 \right) = -\nabla \rho_1 + \frac{1}{Re} \nabla . \vec{\tau_0}$$
$$\frac{\partial \rho_0 Y_0}{\partial t} + \nabla . \left(\rho_0 \vec{v_0} Y_0 \right) = \frac{1}{ReSc} \nabla . \left(\mu \nabla Y_0 \right)$$

• New equation ($\rho_0 = 1/T_0$):

$$\frac{\partial \rho_0}{\partial t} = -\vec{v}_0 \ \nabla \rho_0 - \frac{1}{T_0} \left[\frac{1}{RePr} \nabla \cdot (\mu \ \nabla T_0) \right]$$



Numerical algorithms for Low-Mach

Different methods, e.g. Adams-Bashforth explicit scheme

start by integrating ρ equation:

$$\frac{\rho^{n+1}-\rho^n}{\Delta t} = \left[\frac{3}{2}F^n - \frac{1}{2}F^{n-1}\right], \ \left(\frac{\partial\rho}{\partial t}\right)^{n+1} = \frac{3\rho^{n+1}-4\rho^n+\rho^{n-1}}{2\Delta t}$$

prediction step

$$\frac{\rho^{n+1}\hat{q}_c - \rho^n q_c^n}{\Delta t} = \left[\frac{3}{2}G_c^n - \frac{1}{2}G_c^{n-1} - \mathcal{G}_c\rho^n\right]$$

correction step

$$\mathcal{L}\Phi = \frac{1}{\Delta t} \left[\mathcal{D}\left(\rho^{n+1}\hat{q}_{c}\right) + \left(\frac{\partial\rho}{\partial t}\right)^{n+1} \right],$$

 $(\rho q_c)^{n+1} - \rho^{n+1} \hat{q}_c = -\Delta t \mathcal{G}_c \Phi$



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From fluids to superfluids

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Conclusion

Investigation of the physics of vortex rings

The vortex ring as a fundamental flow

• injection of fluid in a quiescent ambiance.



Renewal of fundamental studies:

Gharib et al., JFM, 1998; Kaplanski et al., Phys. Fluids, 2005.



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Conclusion

Investigation of the physics of vortex rings

Vortex ring simulations

• idea: simulate separately the pipe-flow and the vortex ring





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Conclusion

Investigation of the physics of vortex rings

Postformation evolution: evolution laws

I. Danaila and J. Hélie, <u>Physics of Fluids</u>, 2008. high resolution \rightarrow good agreement with experiments reconciliate Dabiri et al., JFM, 2004 and Maxworthy, JFM, 1972.



Conclusion

Investigation of the physics of vortex rings

Postformation evolution: fit to ideal vortex models

I. Danaila and J. Hélie, Physics of Fluids, 2008.



Conclusion

Improve models for the injection velocity profile

Accurate model for the inflow boundary condition

- I. Danaila, C. Vadean and S. Danaila, Theor. Comput. Fluid Dynamics, 2009.
- Entrance zone + Fully developed zone



• Inlet zone + Stokes zone





N-S Codes

Utility of the DNS of axisymmetric injection flows

Improve models for the injection velocity profile

New inflow velocity model: DNS vs experiments

$$U_{\text{SDV}}(t,r) = U_{\text{CL}}(t) F_{inj}(t) U_b(r,t),$$
$$U_b(r,t) = \frac{1}{2} \left\{ 1 + \tanh\left[\frac{1}{4\Theta(t)} \left(1 - \frac{r}{R_{jet}(t)}\right)\right] \right\}$$





N-S Codes

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Improve models for the injection velocity profile

New inflow velocity model: slug-flow models $\Gamma(t)$

$$\Gamma(t) = \frac{U_0^2 R e_{\rm D}}{32\beta^2} \left[\frac{B(t)(B(t) - \alpha)}{(B(t) - \alpha)^2 + \beta^2} + \frac{\alpha}{\beta} \arctan\left(\frac{B(t)\beta}{\frac{1}{2} - \alpha B(t)}\right) \right]$$



Conclusion

Application in automotive industry

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Impulsively starting flows

Bio-mechanics, synthetic jet actuators, etc. Injection flow in internal combustion engines

• Diesel injectors. • New type of gasoline injectors : low pressure, with swirl, multi-point, piezo actuated.



(courtesy of Continental Automotive France)



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Application in automotive industry

Vortex rings in internal combustion engines: direct Diesel injection

• Jet collapse: as a result of the opposite-sign vortex interactions (dipoles).

• industrial CFD two-phase flow simulation/ experiment (courtesy Institut Français du Pétrole)





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Application in automotive industry

Vortex rings in internal combustion engines: direct Diesel injection

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N-S Codes

Utility of the DNS of axisymmetric injection flows

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New nice mathematical developments

VR problem: mathematical formulation

 $\mathcal{L}\psi$



$$= \frac{\partial}{\partial z} \left(\frac{1}{r} \frac{\partial \psi}{\partial z}\right) + \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial \psi}{\partial r}\right) = \begin{cases} -r\omega_0 f(\psi), \text{ in } \Omega_0 \\ 0, \text{ in } \Pi \setminus \overline{\Omega}_c, \end{cases}$$

$$\psi \text{ and } \nabla \psi \text{ are continuous across } \partial \Omega_c$$

$$\psi = k \text{ on } \partial \Omega_c, \quad \psi = 0 \text{ on } Oz, \end{cases}$$

$$\psi + \frac{1}{2} Wr^2 \to 0 \text{ when } r^2 + z^2 \to \infty.$$

parameters: $W, k, \omega_0, f(r, \psi).$

Fundamental studies (70' and 80')

- J. Norbury, Proc. Camb. Phil. Soc., 1972.
- L. E. Fraenkel & M. S. Berger, Acta Math, 1974.
- H. Berestycki, E. F. Cara & R. Glowinski, RAIRO, 1984.
- C. J. Amick & L. E. Fraenkel, Arch. for Rational Mech. and Analysis, 1987.

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New nice mathematical developments

VR with fixed elliptic $\partial \Omega_b$: non-trivial solutions

Y. Zhang and I. Danaila, Applied Math. Modelling, 2013.





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Conclusion

New nice mathematical developments

Reconstruction of the velocity field



PIV image (Siemens Automotive)



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Difficulties

- the solution is not unique,
- the solution depends on the vortex ring model,
- needs careful set of the matching functional,
- numerics based on non-linear fit procedures.

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Conclusion

New nice mathematical developments

The optimal control problem

Y. Zhang and I. Danaila, J. of Numerical Mathematics, 2012.



$$\begin{split} \min_{\mathbf{X}\in\mathbb{R}^{n}} J(\psi) &= \int_{\partial\Omega_{R}} \left| \frac{1}{r} \left(\frac{\partial\psi}{\partial\vec{n}} - \frac{\partial\psi_{\mathsf{exp}}}{\partial\vec{n}} \right) \right|^{2} \mathrm{d}S \\ \text{subject to} \\ \left\{ \begin{array}{ll} \mathcal{L}\psi &= \omega_{0}f(\mathbf{x},\psi,\mathbf{X}), \quad \text{in }\Omega_{\mathrm{R}}, \\ \psi &= \psi_{\mathsf{exp}}, \qquad \text{on }\partial\Omega_{R}. \end{array} \right. \end{split}$$

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- Find the "best" matching parameters X quasi-Newton Broyden-Fletcher-Goldfarb-Shanno (BFGS) algorithm
- vorticity amplitude (vortex intensity) ω_0

$$\Gamma_{\exp} = \int_{\partial\Omega} \frac{1}{r} \frac{\partial \psi_{\exp}}{\partial \vec{n}} \mathrm{d}\boldsymbol{S} = \omega_0 \int_{\Omega} f(\mathbf{x}, \psi, \mathbf{X}) \mathrm{d}r \mathrm{d}\boldsymbol{z}.$$



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New nice mathematical developments

Original approach for the vortex ring reconstruction

I. Danaila and B. Protasz, submitted, 2014.



$$\begin{split} \min_{f \in H^{1}(\Omega_{R})} J(\psi) &= \int_{\partial \Omega_{R}} \left| \frac{1}{r} \left(\frac{\partial \psi(\mathbf{f})}{\partial \vec{n}} - \frac{\partial \psi_{\exp}}{\partial \vec{n}} \right) \right|^{2} \mathrm{d}S \\ \text{subject to} \\ \begin{cases} \mathcal{L}\psi &= \omega_{0} \mathbf{f}(\psi), \quad \text{in } \Omega_{\mathrm{R}}, \\ \frac{\omega}{r} &= \mathbf{f}(\psi). \end{cases} \end{split}$$

• Optimal problem as in shape optimization weighted Sobolev gradient methods for the minimization Numerical algorithm

- validated against Hill and Norbury vortices,
- used for DNS generated vortex rings.



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Bose-Einstein condensate

New state of the matter: super-atom Properties: superfluid, super-conductor.



Created in 1995 Nobel Prize 2001

C. E. Wieman (Univ. Colorado) E. A. Cornell (Univ. Colorado) W. Ketterle (MIT, Cambridge)



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Conclusion

Quantized vortices: 3D simulation of real experiments (I. Danaila, Phys. Rev. A, 2003, 2004, 2005.)



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Thanks Paolo for the DNS method and code in cylindrical coordinates! They proved useful for many physical and mathematical problems!



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